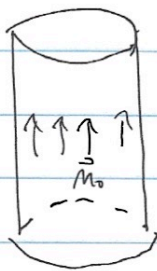


Jackson 5.19(a)

$$\vec{\nabla} \cdot \vec{H} = -\vec{\nabla} \cdot \vec{M}, \quad \vec{\nabla} \times \vec{H} = \vec{J}_f = 0 \Rightarrow \vec{H} = \vec{\nabla} \Phi$$

$$\nabla^2 \Phi = \vec{\nabla} \cdot (\vec{H}) = -\vec{\nabla} \cdot \vec{M}$$

Φ satisfies Poisson's eq with "magnetic polarization charge" $\rho = +\vec{\nabla} \cdot \vec{M}$

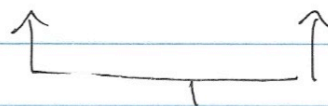


$$\vec{M} = M_0 \hat{z} \Rightarrow \vec{\nabla} \cdot \vec{M} = 0 \text{ inside cylinder.}$$

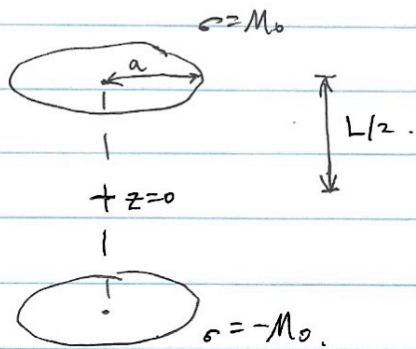
Yet applying Divergence theorem shows that there is $\vec{\nabla} \cdot \vec{M}$ at top and bottom surface.

$$\int_V (\vec{\nabla} \cdot \vec{M}) d\tau = \oint_S \vec{M} \cdot d\vec{a} = 0$$

$$= M_0(\pi a^2) - M_0(\pi a^2) = 0.$$



This implies there is "charge density" M_0 at top surface and $-M_0$ at bottom surface.



$$\Phi = \frac{1}{4\pi} \int \frac{\rho(\vec{x}')}{|\vec{x} - \vec{x}'|} d^3 x'$$

$$= \frac{1}{4\pi} \int \frac{\sigma(\vec{x}')}{|\vec{x} - \vec{x}'|} da_{\text{top}} + \int \frac{\sigma(\vec{x}')}{|\vec{x} - \vec{x}'|} da_{\text{bottom}}$$

$$\Phi = \frac{1}{4\pi} \int \frac{rd\phi dr M_0}{\sqrt{\left(\frac{L}{2} - z\right)^2 + r^2}} - \frac{1}{4\pi} \int \frac{rd\phi dr M_0}{\sqrt{\left(\frac{L}{2} + z\right)^2 + r^2}}$$

$$= \frac{1}{2} \int_0^a \frac{r dr M_0}{\sqrt{\left(\frac{L}{2} - z\right)^2 + r^2}} - \frac{1}{2} \int_0^a \frac{r dr M_0}{\sqrt{\left(\frac{L}{2} + z\right)^2 + r^2}}$$

$$= \frac{M_0}{2} \left\{ \int_0^a \frac{r dr}{\sqrt{\left(\frac{L}{2} - z\right)^2 + r^2}} - \int_0^a \frac{r dr}{\sqrt{\left(\frac{L}{2} + z\right)^2 + r^2}} \right\}$$

$$= \frac{M_0}{2} \left\{ \sqrt{\left(\frac{L}{2} - z\right)^2 + r^2} - \sqrt{\left(\frac{L}{2} + z\right)^2 + r^2} \right\}_0^a$$

$$\Phi = \frac{M_0}{2} \left[\sqrt{\left(\frac{L}{2} - z\right)^2 + a^2} - \sqrt{\left(\frac{L}{2} + z\right)^2 + a^2} + 2z \right]$$

We obtain \vec{H} from $\vec{H} = \vec{\nabla} \Phi$

Φ has dependence on z , so $H_x = 0, H_y = 0,$

$$H_z = \frac{\mu_0}{z} \left[\frac{z \left(\frac{L}{2} - z \right)}{\sqrt{\left(\frac{L}{2} - z \right)^2 + a^2}} \left(\frac{1}{z} \right) (-1) - \frac{z \left(\frac{L}{2} + z \right)}{\sqrt{\left(\frac{L}{2} + z \right)^2 + a^2}} \left(\frac{1}{z} \right) + 2 \right]$$

$$= \frac{\mu_0}{z} \left[\frac{\left(\frac{L}{2} - z \right)}{\sqrt{\left(\frac{L}{2} - z \right)^2 + a^2}} (-1) - \frac{\left(\frac{L}{2} + z \right)}{\sqrt{\left(\frac{L}{2} + z \right)^2 + a^2}} + 2 \right]$$

$$= \frac{\mu_0}{z} \left[2 - \frac{\left(\frac{L}{2} - z \right)}{\sqrt{\left(\frac{L}{2} - z \right)^2 + a^2}} - \frac{\left(\frac{L}{2} + z \right)}{\sqrt{\left(\frac{L}{2} + z \right)^2 + a^2}} \right]$$

From $\vec{H} = \frac{1}{\mu_0} \vec{B} - \vec{M}$, we obtain $\vec{B} = \mu_0 (\vec{H} + \vec{M})$, thus

$$\vec{B} = \begin{cases} \frac{\mu_0}{z} \frac{\mu_0 M_0}{z} \left[2 - \dots - \frac{\left(\frac{L}{2} + z \right)}{\sqrt{\dots}} \right] & \text{outside cylinder.} \\ \frac{\mu_0}{z} \frac{\mu_0 M_0}{z} \left[4 - \dots - \frac{\left(\frac{L}{2} + z \right)}{\sqrt{\dots}} \right] & \text{inside cylinder.} \end{cases}$$

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